Dusty windshield

Radiative magnetohydrodynamics simulations of IR and UV radiation pressure on dusty AGN tori

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Geometrical thickness of the obscuring torus



Dynamical models of the obscuring torus

Warped disk (e.g., Phinney 1989) blocks large solid angle only with severe warps and twists Clumpiness (Krolik & Begelman 1988) needs unusual magnetic field structure to counter inelastic collisions Magnetocentrifugal wind (e.g., Königl & Kartje 1994) needs large-scale magnetic fields Direct magnetic support (Lovelace et al. 1998) needs large-scale magnetic fields Supernovae (e.g., Wada & Norman 2002) needs more energy than observed Stellar ejecta (e.g., Schartmann et al. 2009) ties torus lifetime to starburst Radiative support (e.g., Chan & Krolik 2016, 2017) has not been fully explored with self-consistent simulations

Our recent radiative magnetohydrodynamics simulations

Simulation code



Simulation parameters

Luminosity

0.1 times Eddington

IR opacity

20 times Thomson if below sublimation

UV opacity

Central mass

0.8 solar mass

80 times Thomson if below sublimation

extrapolatable to realistic AGNs

Optical depth

Thomson: 2 infrared: 40 Angular momentum

flat radial profile

genuinely arbitrary

Simulation strategy



Simulation domain



Radiation-driven inflow-outflow



Dust in polar regions of AGNs



Dust in polar regions of AGNs



green: system axis (100 pc)

Kinematics fits expectations

$$\dot{M}_{\text{wind}} \sim \frac{L_{\text{UV}}}{cv_{\infty}} \qquad \frac{L_{\text{kin}}}{L_{\text{UV}}} \sim \frac{v_{\infty}}{c} \qquad v_{\infty}^2 \equiv \frac{GM}{R_{\text{in}}} \frac{L_{\text{UV}}}{L_{\text{E}}} \frac{\kappa_{\text{UV}}}{\kappa_{\text{T}}}$$

Mass loss rate and speed match observed

$$\dot{M}_{\text{wind}} \sim 0.9 \,\text{M}_{\odot} \,\text{yr}^{-1} \times \left(\frac{M}{10^7 \,\text{M}_{\odot}}\right)^{3/4} \left(\frac{L_{\text{UV}}/L_{\text{E}}}{10^{10} \,\text{m}_{\odot}}\right)^{3/4} \\ v_{\infty} \sim 800 \,\text{km} \,\text{s}^{-1} \times \left(\frac{M}{10^7 \,\text{M}_{\odot}}\right)^{1/4} \left(\frac{L_{\text{UV}}/L_{\text{E}}}{0.1}\right)^{3/4}$$

temporal variation: 10%

Radiation-driven outflow explains AGN outflows

- Covering fractions are close to observed type-2 fraction
 - $\begin{array}{c|c} 0.71 \leq C_{\rm IR} \leq 0.73 \\ 0.77 \leq C_{\rm UV} \leq 0.82 \\ 0.78 \leq C_{\rm soft} \leq 0.83 \end{array} \\ \hline \\ \begin{array}{c} \text{due to inflow and outflow;} \\ \text{same for any central mass} \\ 0.15 \leq C_{\rm hard} \leq 0.28 \end{array} \\ \hline \\ \begin{array}{c} \text{due to inflow;} \\ \text{dependent on central mass} \end{array} \\ \end{array}$
- Flat column density distribution agrees with X-ray studies

Radiation-driven outflow explains AGN obscuration



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Radiation-driven outflow explains AGN obscuration

Lessons about torus-scale inflow and outflow

Cartoon of inflow-outflow torus model



Torus as a flow-through system



Constraint from mass

- UV and IR shave off high-latitude dusty gas
- Mass loss rate is 1 M_o yr⁻¹
- But assuming $M = 10^7 \text{ M}_{\odot}$, $L_{\text{UV}}/L_{\text{E}} = 0.1$, $\tau_{\text{T}} = 1$:



 $\begin{array}{ll} \text{Mass} & \sim 2\pi r_{ds}^2 \tau_{\rm T}/\kappa_{\rm T} & \approx 7 \times 10^3 \, {\rm M}_\odot \\ \text{Orbital period} & 2\pi (GM/r_{ds}^3)^{-1/2} \approx 5 \times 10^3 \, {\rm yr} \end{array}$

Mass must be resupplied from galactic scales

Constraint from angular momentum

- 1. Isotropic pressure delivering vertical support also provides radial support
- 2. Direct UV and reprocessed IR provide additional outward radial momentum



Angular momentum must be sub-Keplerian

Constraint from angular momentum

Observational evidence

NGC 4258

Eddington ratio: 0.0002 to 0.02 dynamics: Keplerian rotation

Circinus

Eddington ratio: 0.18 dynamics: Keplerian rotation with outflow

NGC 1068

Eddington ratio: 0.33 to 1.06 dynamics: rotation falls off more slowly than Keplerian

MCP megamasers resembling Keplerian disks

Eddington ratio: 2 × 10⁻³ dynamics: Keplerian at large radii, sub-Keplerian at small radii

Constraint from angular momentum

- Accretion toward inner edge requires low angular momentum
 - Inflow timescale due to stresses is ~ $[\alpha (H/R)^2\Omega]^{-1}$
 - Mass influx at all radii, as determined by inflow timescale, must be comparable to mass outflux

Angular momentum must be low or rapidly removed

Constraint from energy

- Radiation does positive work on outflows
- Binding energy of torus decreases
- Torus eventually becomes unbound



Constraints on inflow of steady-state torus

- 1. Mass must be resupplied from galactic scales
- 2. Angular momentum must be sub-Keplerian
- 3. Energy must be kept low

How can mass resupply satisfy constraints 2 and 3?

- Stresses in inflow rapidly remove angular momentum and energy
- Mass resupply has inherently low angular momentum and energy

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Summary

- Torus in RMHD simulations settles into steady inflow–outflow
- IR in central hole drives high-latitude, wide-angle outflow with expected:
 - kinematics
 - obscuration properties
- Steady-state irradiated tori must:
 - be resupplied with mass
 - have sub-Keplerian rotation



References I

Asmus, D., Hönig, S. F., & Gandhi, P. 2016, ApJ, 822, 109

- Braatz, J. A., Wilson, A. S., Gezari, D. Y., Varosi, F., & Beichman, C. A. 1993, ApJL, 409, L5
- Chan, C.-H. & Krolik, J. H. 2016, ApJ, 825, 67
- Chan, C.-H. & Krolik, J. H. 2017, ApJ, 843, 58
- Elvis, M., Wilkes, B. J., McDowell, J. C., et al. 1994, ApJS, 95, 1
- Gnedin, N. Y. & Abel, T. 2001, NewA, 6, 437
- Hönig, S. F., Kishimoto, M., Antonucci, R., et al. 2012, ApJ, 755, 149
- Jiang, Y.-F., Stone, J. M., & Davis, S. W. 2014, ApJS, 213, 7
- Königl, A. & Kartje, J. F. 1994, ApJ, 434, 446
- Krolik, J. H. 2007, ApJ, 661, 52
- Krolik, J. H. & Begelman, M. C. 1988, ApJ, 329, 702
- Lovelace, R. V. E., Romanova, M. M., & Biermann, P. L. 1998, A&A, 338, 856
- Nenkova, M., Ivezić, Ž., & Elitzur, M. 2002, ApJL, 570, L9

References II

- Phinney, E. S. 1989, in NATO Advanced Science Institutes Series C: Mathematical and Physical Science, 290, Theory of Accretion Disks, ed.
 F. Meyer, W. J. Duschl, J. Frank, & E. Meyer-Hofmeister (Dordrecht: Kluwer), 457
- Pier, E. A. & Krolik, J. H. 1992, ApJ, 401, 99
- Schartmann, M., Meisenheimer, K., Klahr, H., et al. 2009, MNRAS, 393, 759
- Stone, J. M., Gardiner, T. A., Teuben, P., Hawley, J. F., & Simon, J. B. 2008, ApJS, 178, 137
- Tristram, K. R. W., Burtscher, L., Jaffe, W., et al. 2014, A&A, 563, A82
- Wada, K. & Norman, C. A. 2002, ApJL, 566, L21

Modeling AGN spectral energy distributions



Phenomenological models of the obscuring torus





(Nenkova et al. 2002)



Challenges of simulating irradiated tori



Poloidal and midplane slices of simulation



